



Project X Based Muon Factory

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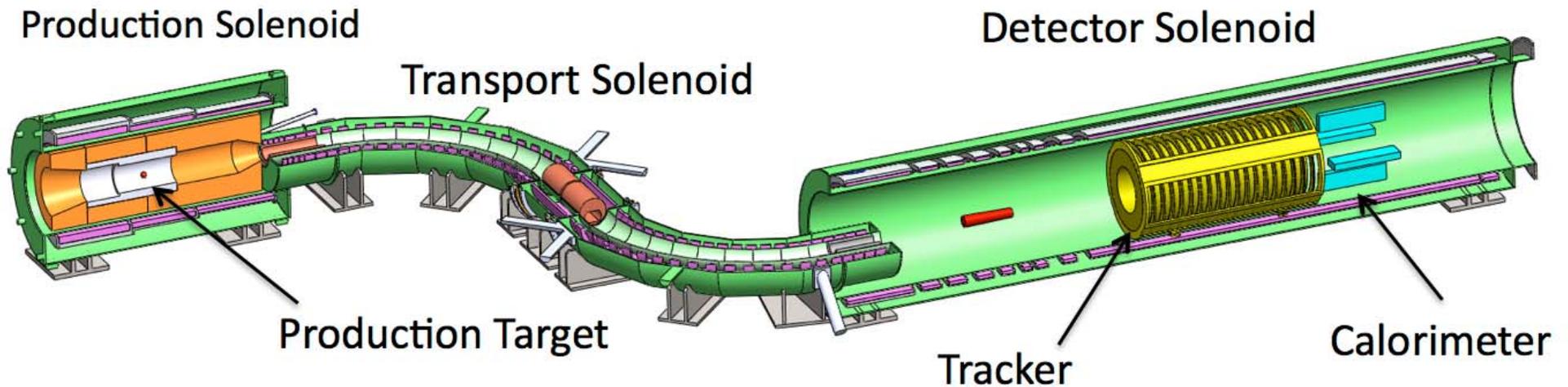
APT Seminar

Fermilab

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Present μ -to- e

- Conversion - $2.1 \cdot 10^{-3}$ ($dN_p/dt = 2.4 \cdot 10^{13} \text{ s}^{-1}$, $P = 25 \text{ kW}$, $dN_\mu/dt = 5 \cdot 10^{10} \text{ s}^{-1}$)
- Extinction $< 10^{-10}$ (sensitivity $6 \cdot 10^{-17}$ (90% C.L.))
- Target (gold, $L \sim 16 \text{ cm}$, $r = 0.5 \text{ cm}$, water cooled)
 - ◆ Total power - 25 kW
 - ◆ Power left in the target - 2 kW
- Secondary target
 - ◆ 17 Al discs, 0.2 mm thick, 5 cm apart, tapered radii - $r_d = 8.3 \rightarrow 6.53 \text{ cm}$
- Magnetic fields
 - ◆ Production solenoid: 5T \rightarrow 2.5 T, internal radius 0.75 m (reflection of muons)
 - ◆ Transport solenoid - 2 T
 - ◆ Detector solenoid : 2T \rightarrow 1T (reflection of electrons with negative $p_{||}$)



Major Requirements to a New Generation μ -to- e Experiment[†]

- ~100 times better than μ -to- e
 - ◆ single event sensitivity $2 \cdot 10^{-19}$ (or $6 \cdot 10^{-19}$ at 90% CL)
 - ⇒ $5 \cdot 10^{18}$ muons: 2 years of $2 \cdot 10^7$ s each
 - ⇒ $1.2 \cdot 10^{11}$ muons/s or 2.5 times of μ -to- e
 - ⇒ Conversion $\geq 6 \cdot 10^{-5}$ (3 GeV beam with 1 MW power)
 - ◆ $P_c < 20$ MeV i.e. $E_{kin} < 1.9$ MeV (stopped in 0.4 mm Al foil)
 - ◆ Extinction $< 10^{-14}$ for pions; no antiprotons
 - ◆ Short pulse: $t < 10$ ns
 - ◆ Detector is located underground (≥ 12 m)
- Short pulse and very good extinction imply that the beam transport has to be in an isochronous beam line
 - ◆ Drastic reduction of transverse and longitudinal acceptances
 - ⇒ 1 MW Project X power should be helpful
- Limitation of maximum energy to < 1 MeV points out to the muon deceleration as a possible choice

[†] Bernstein & Prebys, July 26, 2011

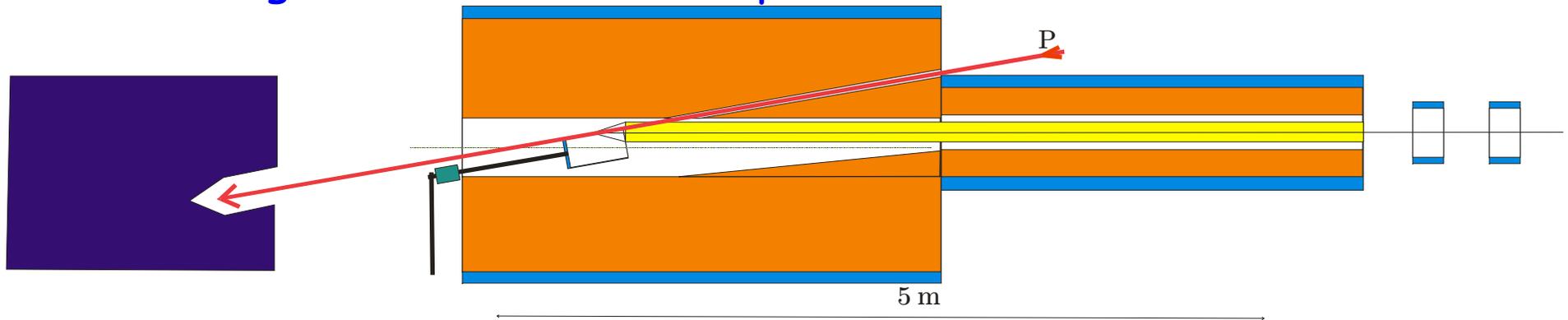
Questions to be addressed

- Target or how to handle the beam power
 - ◆ Required shielding
 - ◆ Target shape and material
 - ◆ Target cooling
- Particle distributions and emittances of pions and muons
 - ◆ Effect of magnetic field strength and line acceptance
 - ◆ Distribution changes at conversion of pions to muons
- Phase space manipulation for muon beam
 - ◆ Which area of the phase space we would like to use
 - ◆ Matching between the target and the decay solenoid
 - ◆ Deceleration and ionization cooling
- Beam transport and background suppression
 - Beam transport in isochronous beam line
 - Deceleration versus ionization cooling

- *This talk should not be considered as a proposal for a project
It rather discusses major limitations and possible ways of their mitigation*

Target

- The target length should be ~ 1.5 of nuclear interaction length
 - \Rightarrow Carbon ~ 60 cm
 - \Rightarrow Tantalum ~ 15 cm
- The beam leaves $\sim 10\%$ of its energy in the target;
 - $\Rightarrow \sim 100$ kW for 1 MW power
- ◆ 90% goes to the beam dump



- Relative to pulsed beam the CW beam drastically reduces stress in target

Target cooling

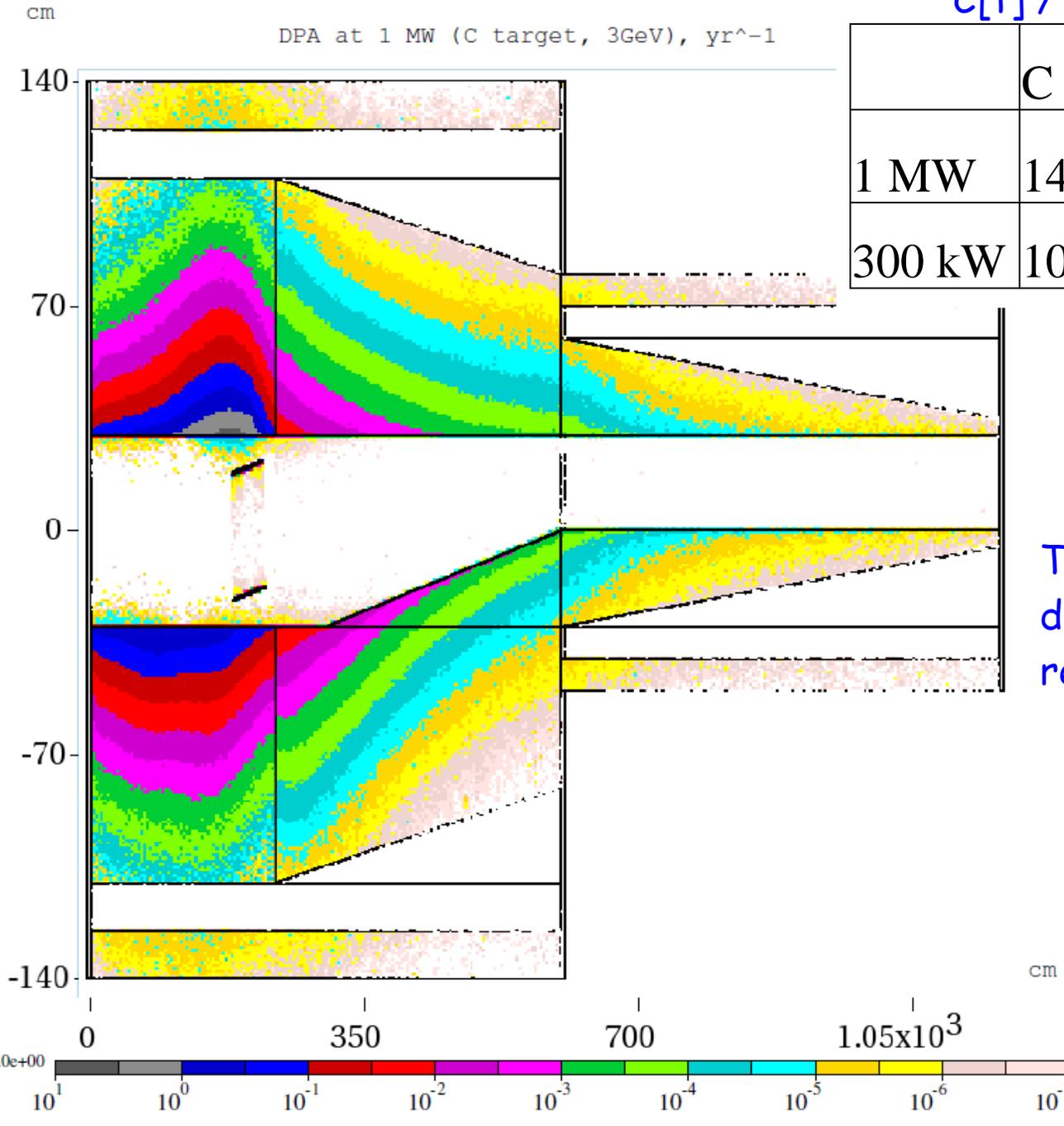
- For 1 MW beam power the power left in the target is ~ 100 kW
 - ◆ Heat cannot be removed from pencil target: $dP/dS \sim 2$ kW/cm² for $R \sim 0.5$ cm
 - ◆ Relative to this an oxidation and repairs look as an easy problem
- Two possibilities
 - ◆ Liquid metal stream (muon collider)
 - Looks expensive
 - Reliability, safety and repair issues
 - ◆ Rotating cylinder cooled by black body radiation
 - PSI uses a rotating graphite target at 1 MW beam power
 - Tantalum, $R=10$ cm, $d=0.5$ cm, $L=15$ cm, 400 rev/min
 - $\Rightarrow T \approx 3000$ K (melting $T = 3270$ K), $\Delta T \approx 50$ C
 - Graphite (C), $R=10$ cm, $d=0.5$ cm, $L=40$ cm, 60 rev/min
 - $\Rightarrow T \approx 1800$ K (melting $T = 3270$ K), $\Delta T \approx 50$ C
 - For C temp. looks OK but we still have to address
 - \Rightarrow Bearing lifetime under radiation (rotation)
- Any solution requires vacuum windows to separate target from the beam \Rightarrow 1 MW windows
 - Do we need to have the target in vacuum?

Effects of radiation

Shielding estimate

$C[t] / W[t] / R_{max} [cm]$

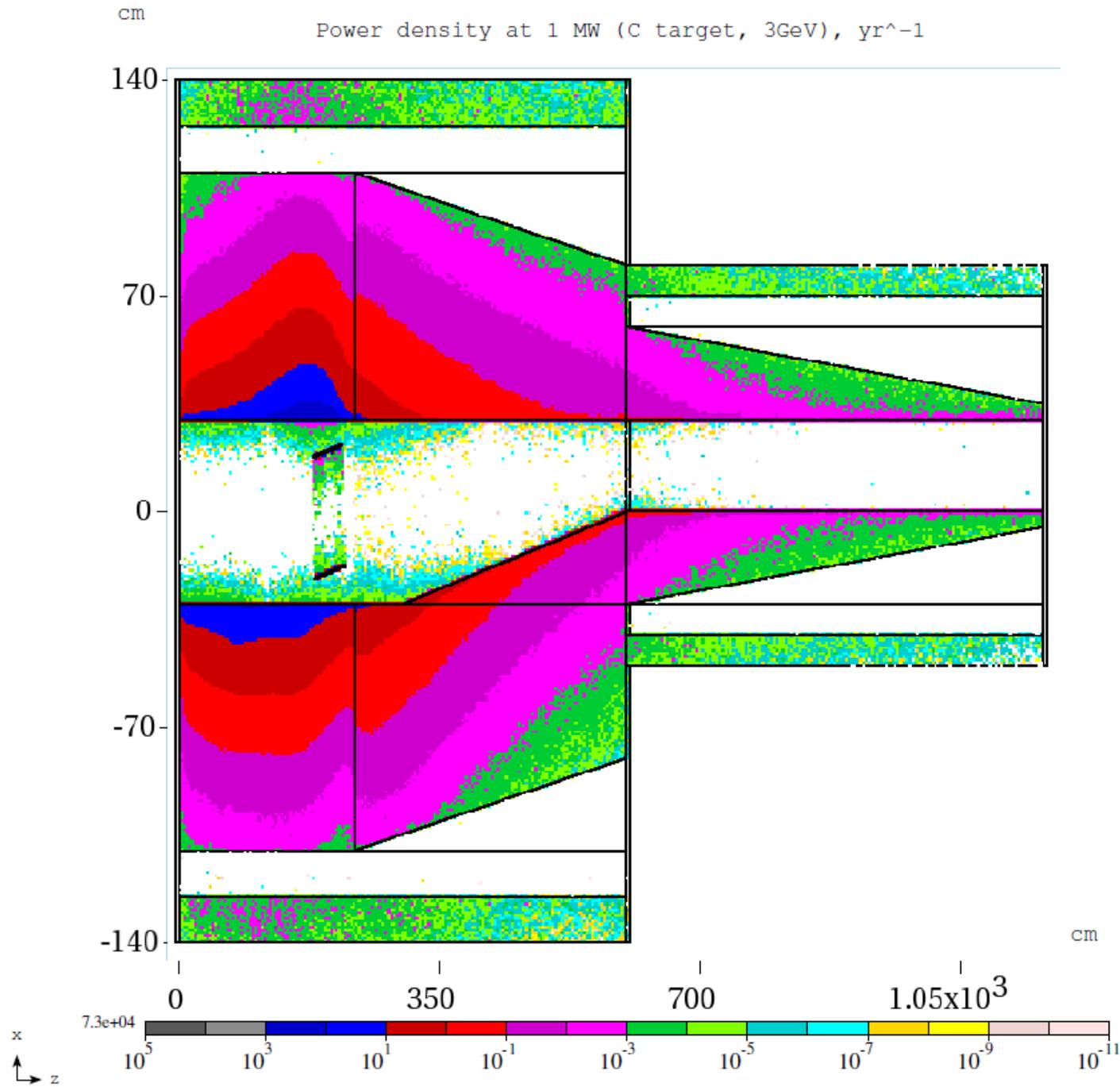
	C target	Ta target
1 MW	140/80 (110)	180/100 (125)
300 kW	100/55 (95)	110/65 (100)



This preliminary absorber design satisfies typical requirements for SC coils

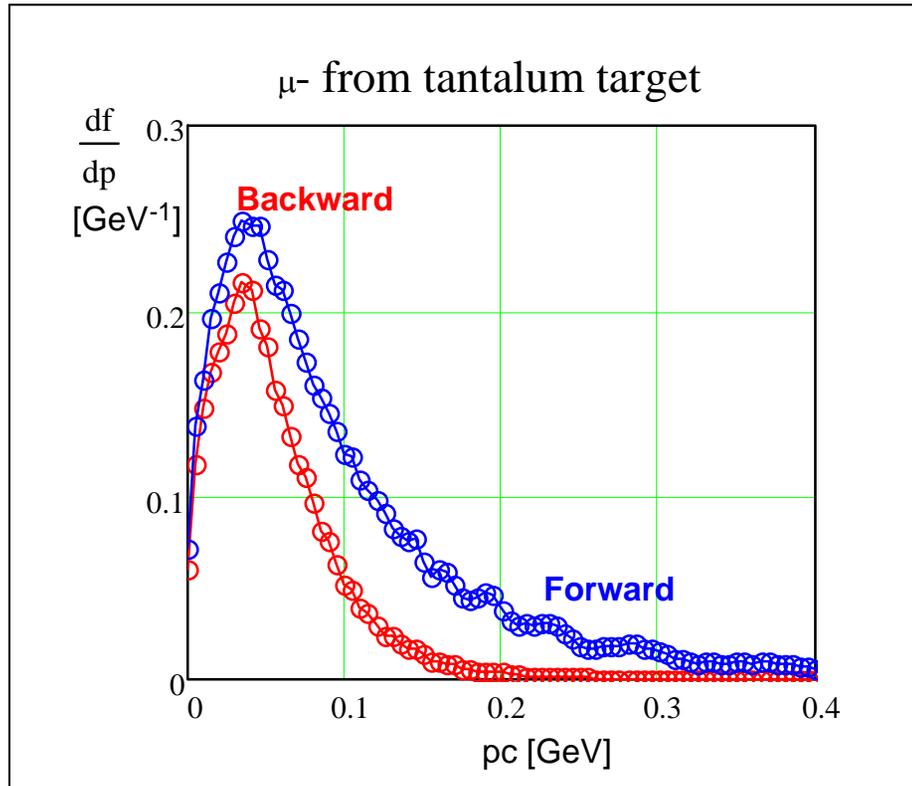
- peak DPA $1E-5 \text{ yr}^{-1}$
- power density $3 \mu\text{W/g}$
- absorbed dose 60 kGy/yr
- Dynamic heat load is 10 W

Effects of radiation (continue)

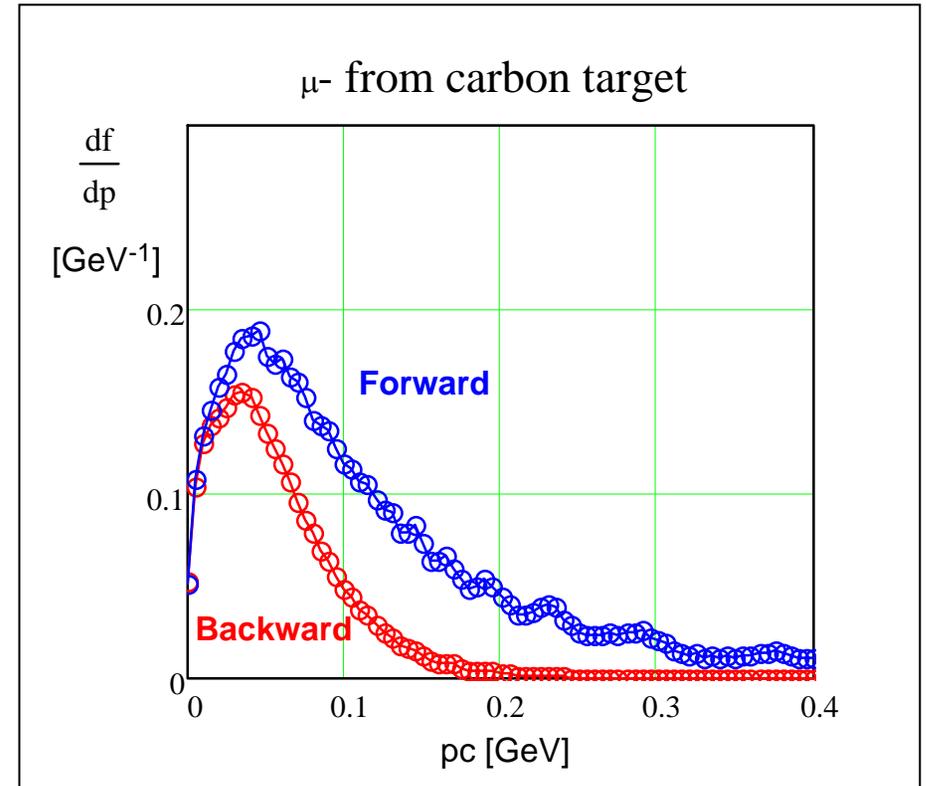


Muon's longitudinal distribution (per 1 GeV of proton energy)

- 3 GeV/c ($E_{kin}=2.2$ GeV) proton beam (this choice is supported by measurements)
 - ◆ $\sigma_x = \sigma_y = 1$ mm - parallel beam, proton multiple scattering unaccounted



*Tantalum hollow cylinder ($P_c=3$ GeV)
 $R_{out}=20$ cm, $\Delta R=5$ mm, $L=16$ cm, $\theta=300$ mrad
Total muon yield at ± 10 m
Forward - 1.4% per proton GeV
Backward - 0.73% per proton GeV*



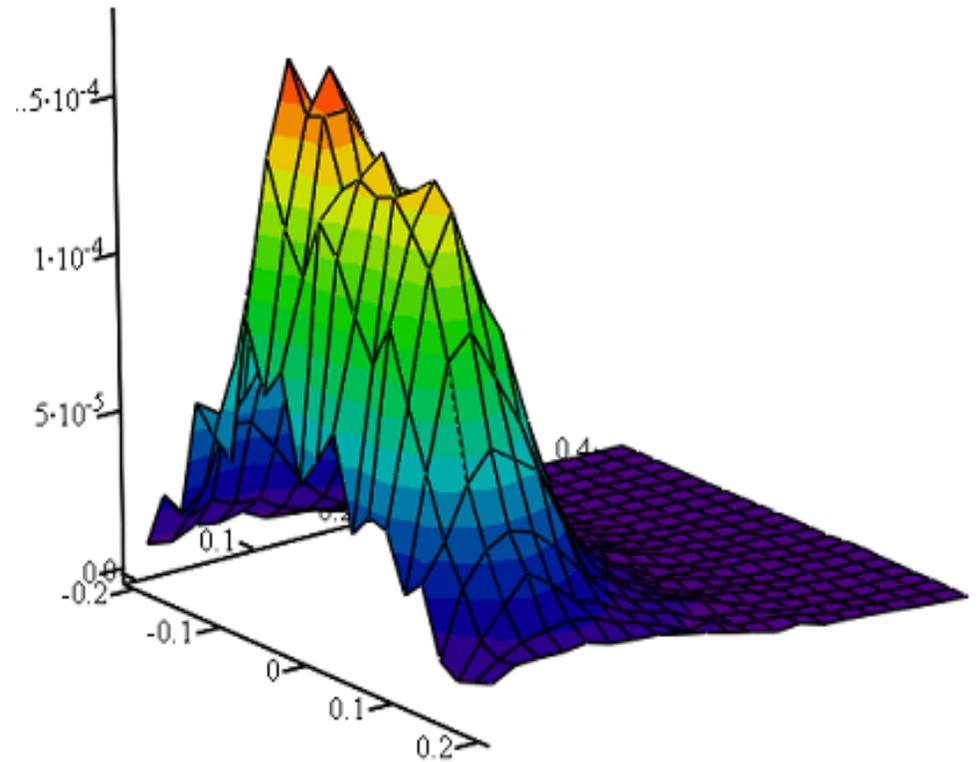
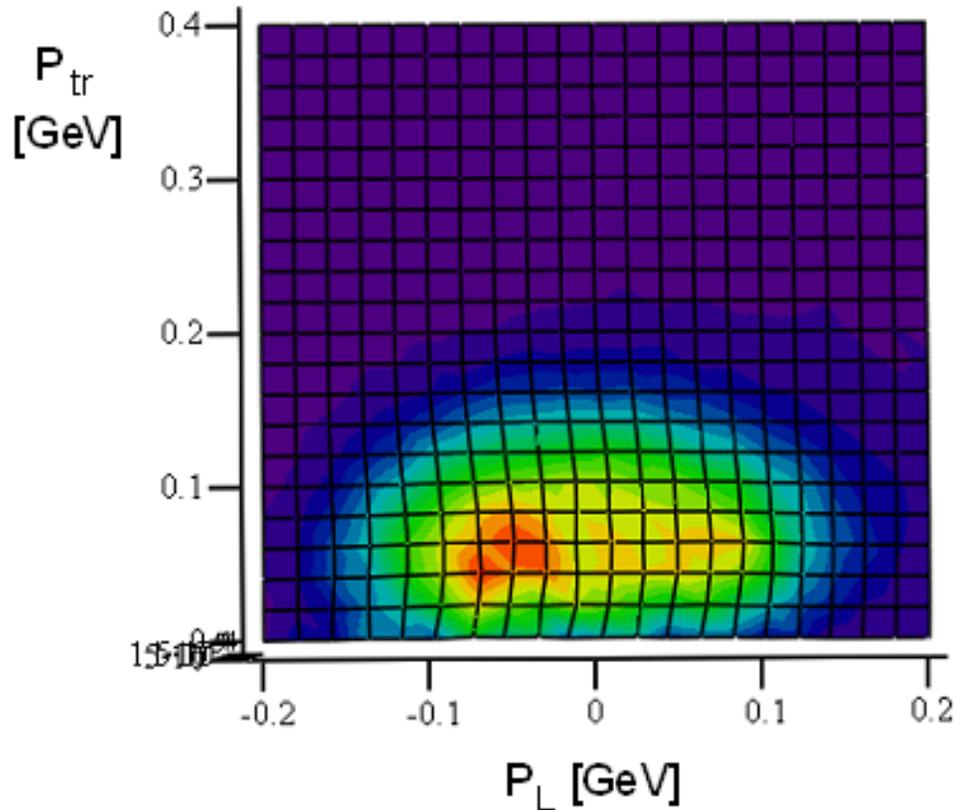
*Carbon hollow cylinder ($P_c=3$ GeV)
 $R_{out}=20$ cm, $\Delta R=5$ mm, $L=40$ cm, $\theta=200$ mrad
Total muon yield at ± 10 m
Forward - 1.3% per proton GeV
Backward - 0.59% per proton GeV*

- Small difference between forward and backward muons for $P_c < 50$ MeV

Muon's longitudinal distribution (contunue)

- Compared to a pencil like target a hollow cylinder target has smaller muon yield by more than factor of 2
 - ◆ But it allows one to use much larger beam power
- For $P_c < 100$ MeV the carbon target has smaller yield but
 - ◆ Less problems with cooling due to larger length
 - ◆ It also makes less neutrons
- Beam damp inside solenoid would be a formidable problem therefore below we assume:
 - ◆ Backward muons
 - ◆ Carbon target
- We also assume the proton energy of 2.21 GeV
 - For $E_{kin} \in [2, 8]$ the production of slow muons per unit beam power weakly depends on the beam energy

Pion distribution over momentum



Pion distribution over momentum, d^3N/dp^3 ,

Nickel cylinder, $L=10$ cm, $r=0.4$ cm; no magnetic field

Total production: forward 5.3% $p_{\perp} \text{ GeV}^{-1}$; backward - 2.9% $p_{\perp} \text{ GeV}^{-1}$

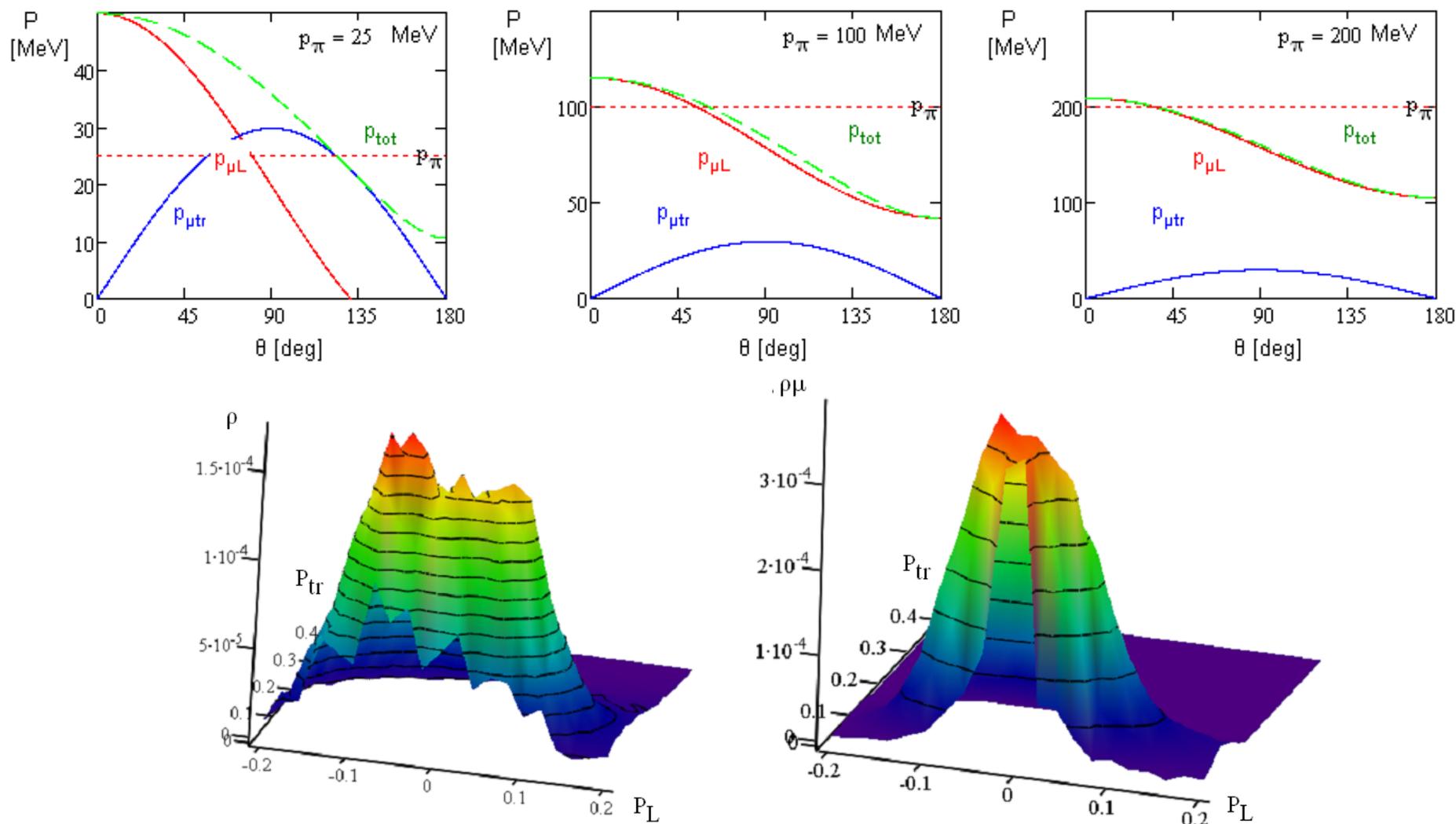
- Pion distribution is close to the Gaussian one with $\sigma_p = 100$ MeV/c
- Distribution function approaches zero due to particle deceleration at the

target surface: $f(p_{\perp}) \propto p^3 / (p^4 + p_r^4)^{3/4}$, $p_r \approx \sqrt[4]{m_{\mu}^3 \delta r (dE/dx)_0 / c}$

where $(dE/dx)_0$ is the energy loss at $\beta = 1$ (~ 1.6 MeV/(g/cm²)); $p_r \approx 20$ MeV/c

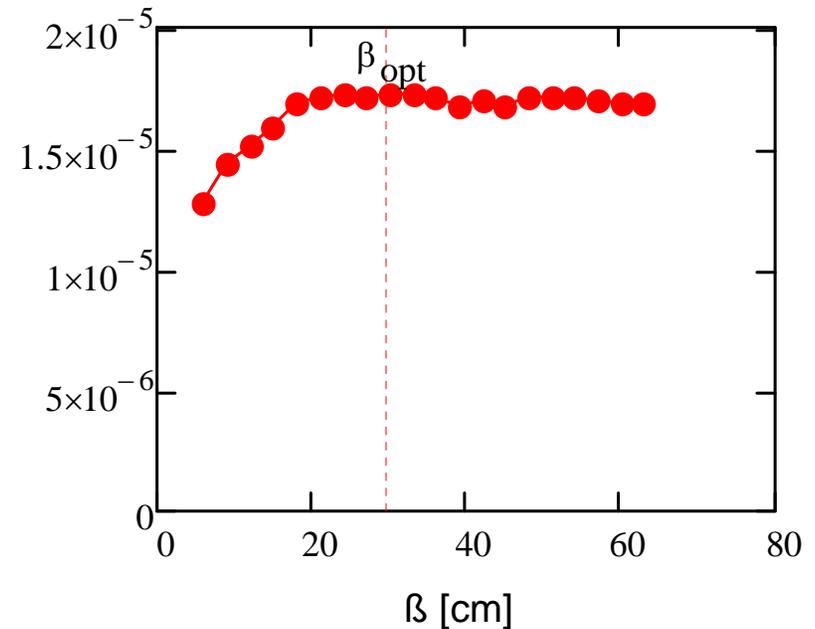
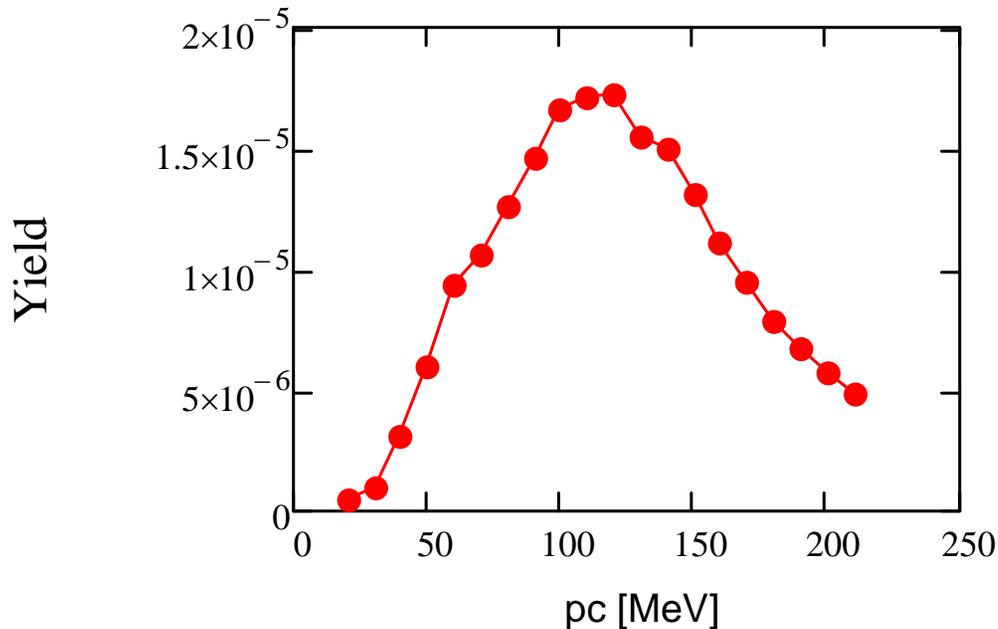
Muon distribution over momentum

- After decay a muon inherits the original pion momentum with Δp correction depending on the angle of outgoing neutrino, $\Delta p_{cm} = 29.8 \text{ MeV}/c$
- For most of pions ($p > 60 \text{ MeV}/c$) a decay makes a muon with smaller p
 \Rightarrow Momentum spread in μ -beam is smaller than in π -beam



Muon yield into the beamline finite acceptance

- Issues related to the beam transport in a “normal” beam line
 - ◆ Isochronous transport requires bends and soft focusing, $Q_x \sim 1$
 - ◆ Beam line limits maximum acceptance and momentum spread to $\varepsilon \approx 0.3\text{-}3\text{ cm}$, $\Delta p/p \approx \pm 0.15$
 - ◆ We can match the beam line to the decay solenoid so that to maximize the capture $\Rightarrow \beta_{\text{opt}}$
 - Weak dependence on β for small ε
 - Strong suppression of small energy muons ($pc < 50\text{ MeV}$) due to deceleration in medium

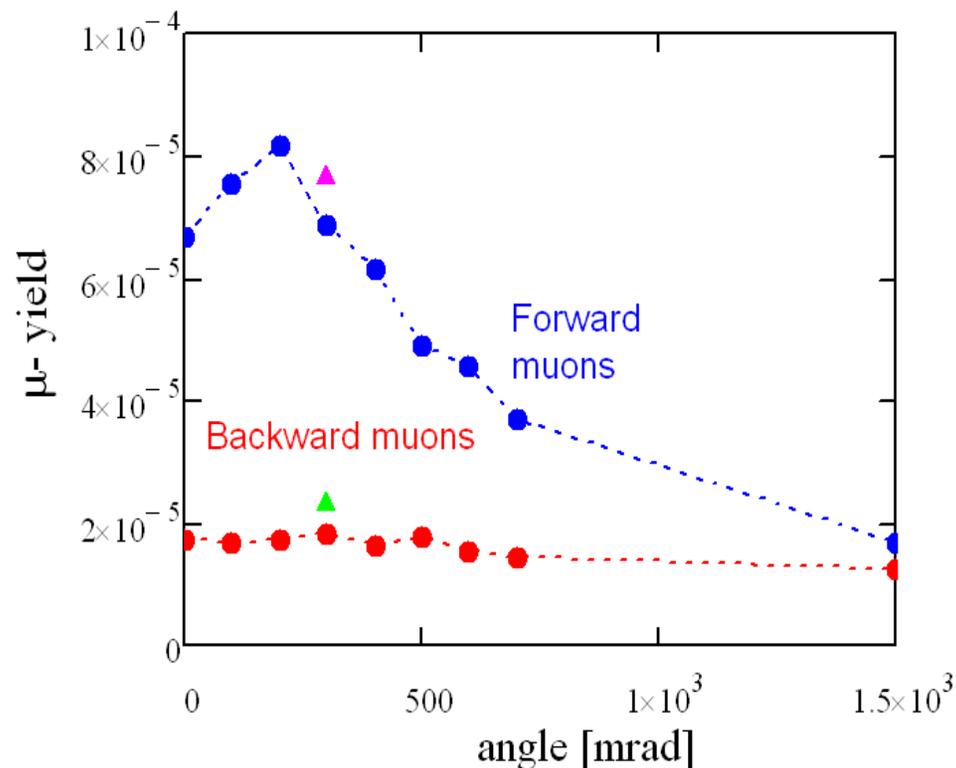
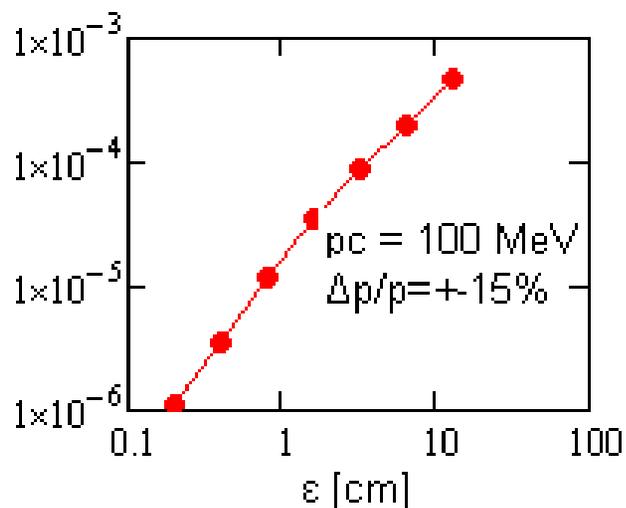


Graphite target, backward muons, $\varepsilon_x = \varepsilon_y = 1\text{ cm}$, $\Delta p/p = \pm 0.15$, $\theta = 200\text{ mrad}$, $B = 2.5\text{ T}$.

Muon yield into the beamline finite acceptance (continue)

- Pions are produced with zero angular momentum and muons inherit this property
 - ⇒ Absence of x-y correlations after beam exit from magnetic field
 - ⇒ Beam line axis has to coincide with solenoid axis
- Yield is proportional to B_{target}
 - ◆ 2.5 T \rightarrow 5 T would double the yield
- Yield is $\propto \Delta p/p$ (for $\Delta p/p \ll 1$)
- Yield is $\propto \varepsilon^{1.5}$

Yield, C cylinder, backward μ

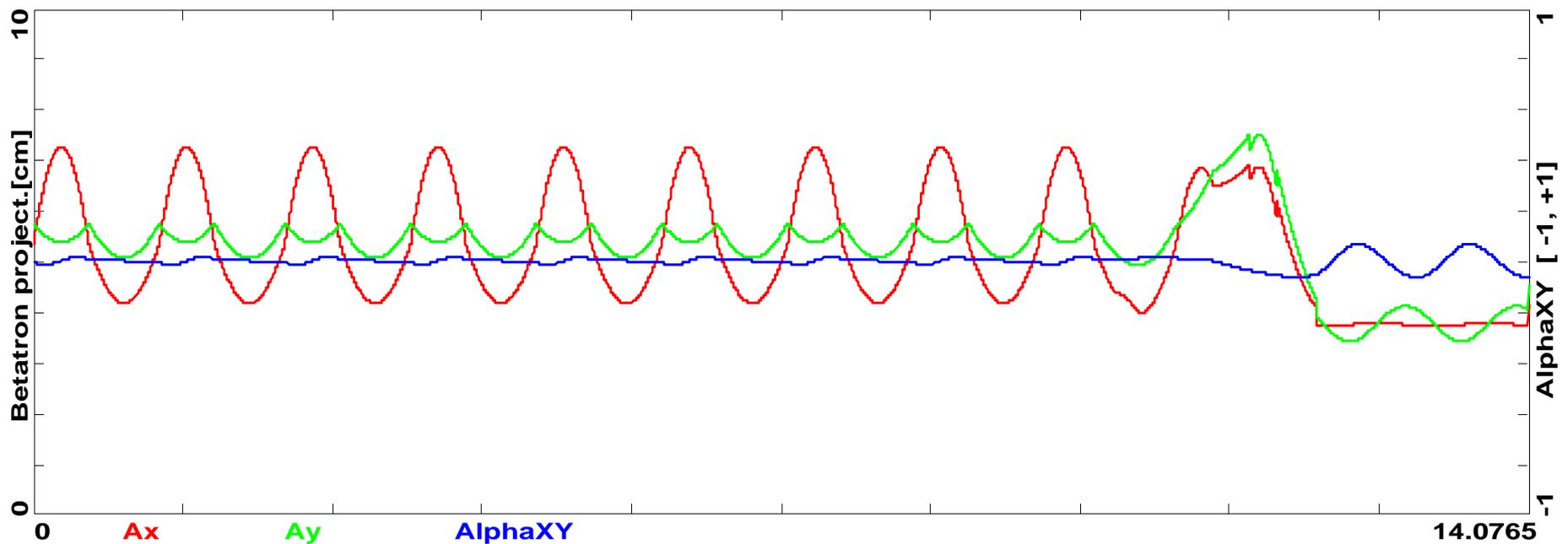


*Dependence of muon yield on target angle relative to magnetic field for carbon target into the following phase space: $\varepsilon_x = \varepsilon_y = 1$ cm, $\Delta p/p = \pm 15\%$, Optimal momenta are: 100 MeV/c for backward and 200 MeV/c for forward muons
Triangles show results for tantalum target*

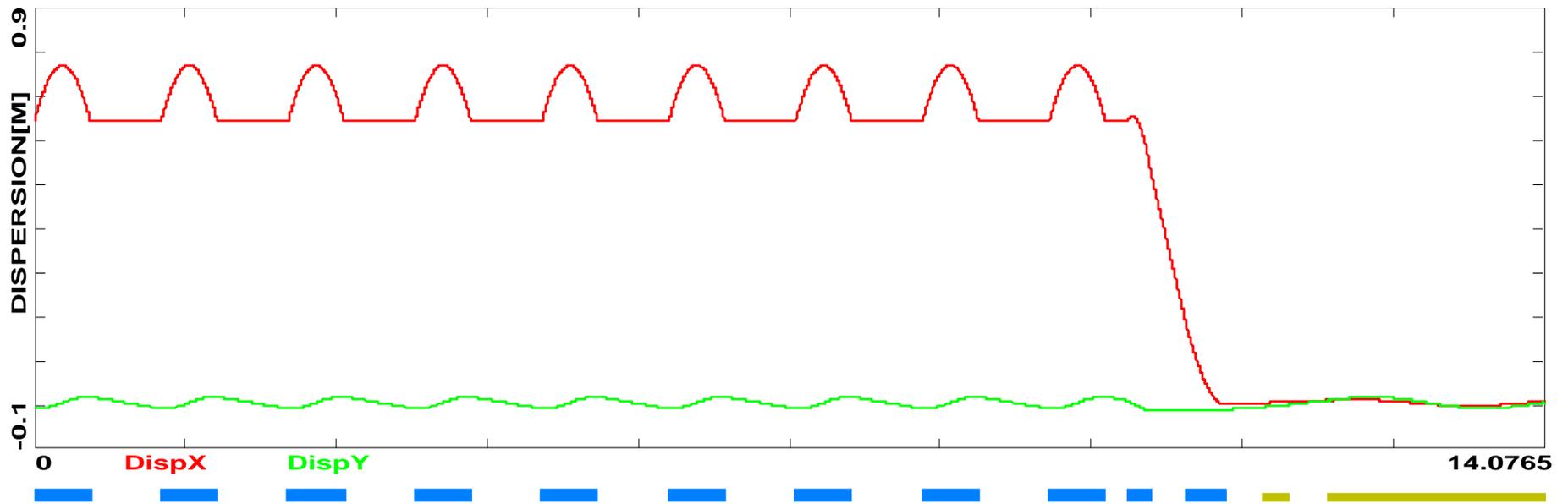
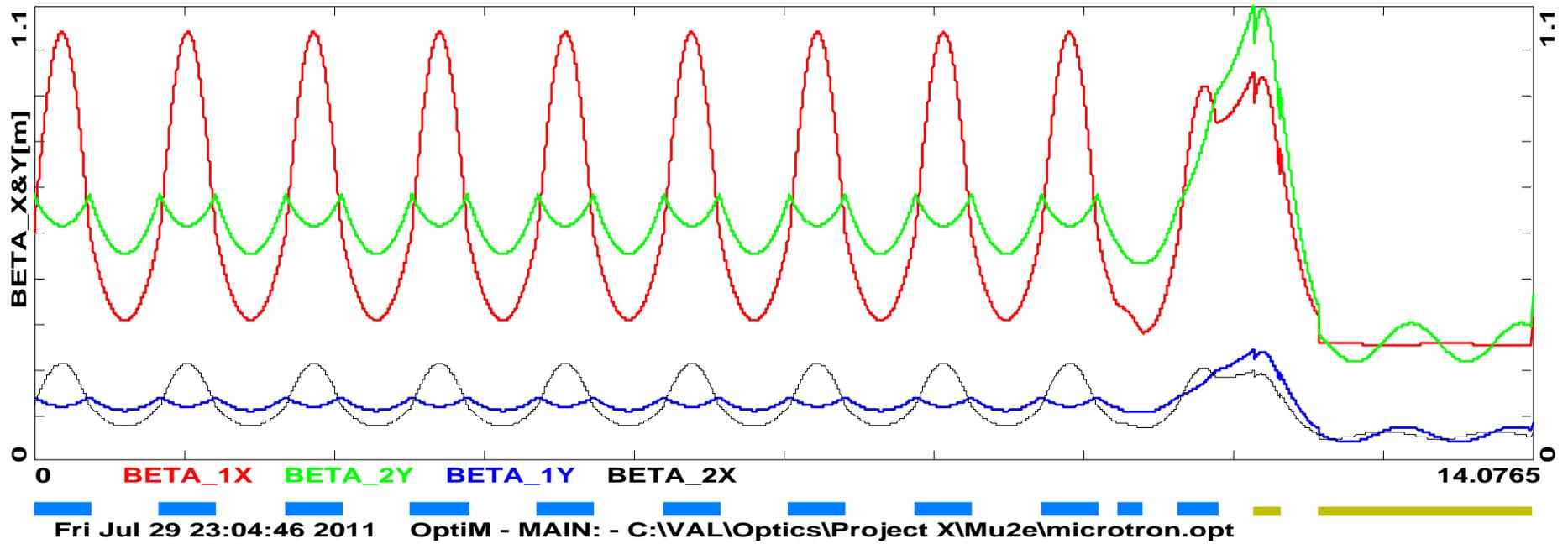
Beam transport in Helical Transport Line

- The line consists of downward spiral
 - ◆ matched to the production and detector solenoids with two dipoles and one or two solenoids at each end
- One revolution includes 4 dipole magnets: $B=5$ kG ($P_c=50$ MeV), $L=52.3$ cm, $R=33.3$ cm, gap 13 cm, good field region width: ± 15 cm
- The line acceptance 0.41 cm; Momentum spread ± 0.15 , it descends with angle of 2.591 deg, step of the helix is 23.973 cm

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Betatron beam envelopes for helix and match to the detector solenoid. Acceptance 0.41 cm



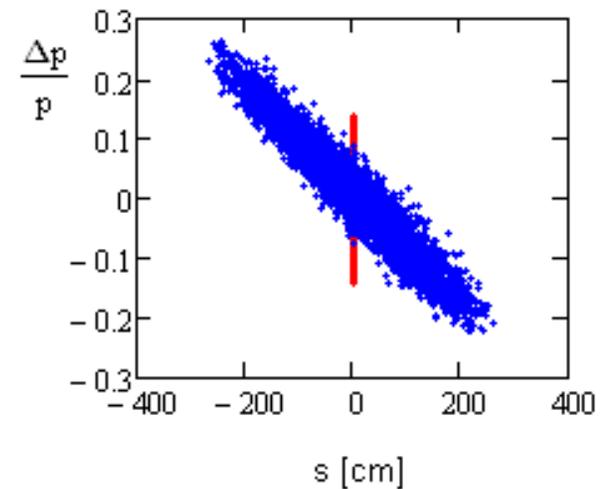
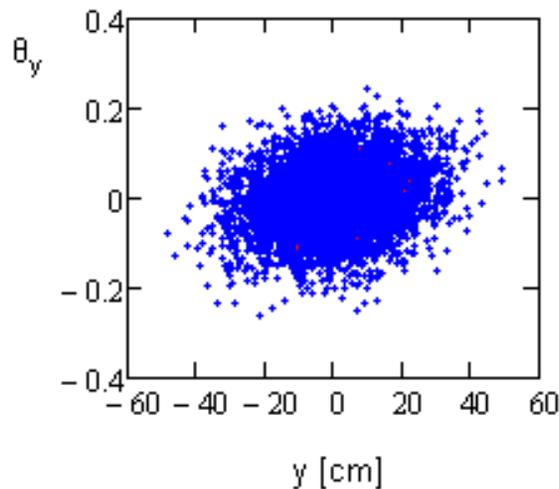
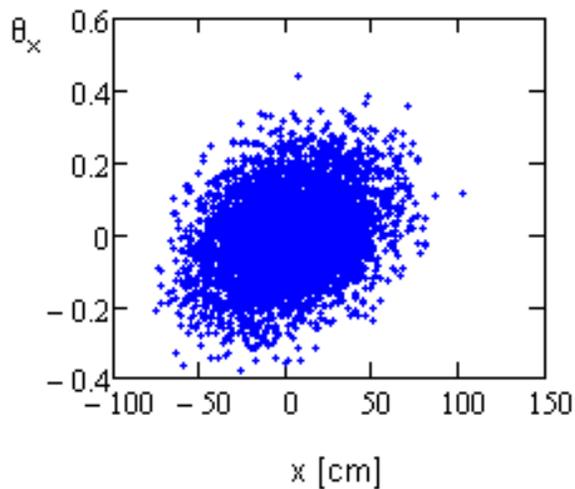
4D beta-functions (top) and dispersions (bottom) for helix and match to the detector solenoid

Beam transport limitations

- To match the yield requirement of $\sim 10^{-4}$ we need to have a line with acceptance of ~ 3 cm (backward muons from carbon target)
 - ◆ Similarity of optics yields: $\varepsilon \propto a \propto \beta_{x,y} \propto R_0$
 - ◆ Isochronicity requires soft focusing, $Q_x \sim 1$
 - ◆ Magnetic fields are reduced with increase of R_0 making magnet price affordable
 - ◆ Total length and number of turns is determined by required pion extinction (~ 70 m for 50 MeV/c and extinction of 10^{-14})

Possibilities with Deceleration

- Deceleration in electro-magnetic structure results in the adiabatic antidumping, with consequential 6D emittance growth $\propto p^{-3}$, i.e. 8 times for every factor of 2 in momentum
- Deceleration in the material looks much better at large p ($p \geq m_\mu$) but behaves the same way ($\propto p^{-3}$) for non-relativistic particles
 - ◆ even worse than it if multiple scattering is important (large $\beta_{x,y}$ at absorber)
 - ◆ Redistribution of damping decrements in realistic simulation partially helps but does not address the problem



$g_L \equiv 1$	$\mu_x \equiv 2 \cdot \pi \cdot 0.25$	$\beta_x \equiv 200$ cm	$\alpha_x \equiv -0.3$
$\kappa_{\text{scat}} \equiv 1$	$\mu_y \equiv 2 \cdot \pi \cdot 0.25$	$\beta_y \equiv 200$ cm	$\alpha_y \equiv -0.2$
	$D \equiv 150$ cm	$D_p \equiv 0.0$	$M_{56} \equiv 0$
	$\varepsilon_x \equiv 3$ cm	$\varepsilon_y \equiv 3$ cm	$\sigma_p \equiv 0.15$

$\kappa_{\text{eff}} = 0.281$

$\frac{\varepsilon_x^{\text{fin}}}{\varepsilon_x^{\text{in}}} = 6.89$

$\frac{\varepsilon_y^{\text{fin}}}{\varepsilon_y^{\text{in}}} = 2.54$

$\frac{\sigma_p^{\text{fin}}}{\sigma_p^{\text{in}}} = 1.758$

Conclusions

- Requirements on the total muon flux accepted into a beam line can be met for muons with momenta ~ 100 MeV ($E_{\text{kin}}=40$ MeV)
- Requirements on isochronicity can be met
- Number of muons at low energy is reducing fast

- ◆ Deceleration results in about the same yield decrease as the direct capture would do
- ◆ Changes in the target design look as a promising avenue

$$f(p_{\perp}) \propto p^3 / (p^4 + p_r^4)^{3/4}$$

$$p_r \approx \sqrt[4]{m_{\mu}^3 \delta r (dE/dx)_0 / c}$$

- ◆ Small emittance of Project X will be helpful
 - Convergent beam
 - Mitigation of multiple scattering for protons in the target

